

Problem 1. Let $\psi : \mathbb{R} \rightarrow \mathbb{R}$ be the $C_c^\infty(\mathbb{R})$ function with $\text{supp } \psi = [-1, 1]$ (this function was constructed on page 236 of Folland, but you don't need its explicit form, just its properties).

- (a) Define the sequence $\{f_k\}_{k \in \mathbb{N}} \subset C_c^\infty(\mathbb{R})$ by $f_k(x) = \frac{1}{k} \cos\left(\frac{k\pi x}{2}\right) \psi(x)$. Explain why sequence does *not* converge to 0 in $\mathcal{D}(\mathbb{R})$, although pointwise $f_k \rightarrow 0$.
- (b) Let $\{g_k\}_{k \in \mathbb{N}} \subset C_c^\infty(\mathbb{R})$ be a sequence defined by $g_k(x) = e^{-k} \psi(x - k)$. Does this sequence converge in the uniform metric? Does it converge in $\mathcal{D}(\mathbb{R})$? Explain briefly.

Problem 2. For $t > 0$ define the functions

$$f_t : \mathbb{R} \rightarrow \mathbb{R} : x \mapsto f_t(x) = \frac{1}{t} \chi_{[0,t]}(x) .$$

Clearly, each f_t is in $L^1_{\text{loc}}(\mathbb{R})$, and, thus, defines a distribution $f_t \in \mathcal{D}'(\mathbb{R})$ by

$$\langle f_t, \phi \rangle = \int_{\mathbb{R}} f_t(x) \phi(x) dx , \quad \phi \in \mathcal{D}(\mathbb{R}) .$$

- (a) Find $\lim_{t \rightarrow 0+} \langle f_t, \phi \rangle$. Can you write your result as $\int \phi d\mu$ for some measure μ ?
- (b) If τ_s stands for the translation operator (see page 238 of Folland), find $\langle \tau_s f_t, \phi \rangle$ and the limits $\lim_{t \rightarrow 0+} \langle \tau_s f_t, \phi \rangle$ and $\lim_{s \rightarrow 0} \lim_{t \rightarrow 0+} \left\langle \frac{\tau_s f_t - f_t}{s}, \phi \right\rangle$. Can you write your results in the form $\int \phi d\mu$ for some measures μ ?

Problem 3. Let $X = [0, 1]$ and \mathcal{M} be the σ -algebra of Borel subsets of X . Let $F(t) = t^2$, $G(t) = t^3$, and define the measures ϕ and γ on \mathcal{M} by

$$\phi(E) = \int_E 1 d\mu_F , \quad \gamma(E) = \int_E 1 d\mu_G , \quad E \in \mathcal{M} .$$

Does $\frac{d\phi}{d\gamma}$ exist? Does $\frac{d\gamma}{d\phi}$ exist? Compute the values of the Radon-Nikodym derivatives that exist. Justify.

Problem 4. Consider the *heat equation* in one time and one space dimension,

$$\partial_t u(t, x) = \partial_{xx} u(t, x) , \quad t > 0 , \quad x \in \mathbb{R}$$

(the name comes from the fact that it describes the propagation of heat in a homogeneous medium without heat sources in one spatial dimension; here $u(t, x)$ is the temperature of the medium at time t at the location with spatial coordinate x). Usually one imposes an initial condition, i.e., a function u_0 such that $\lim_{t \rightarrow 0+} u(t, x) = u_0(x)$, and sometimes also conditions on the behavior of the solution as $x \rightarrow \pm\infty$.

- (a) Let for $t > 0$ and $x \in \mathbb{R}$, the function $u(t, \cdot)$ be a solution of the heat equation at time t that tends to zero fast enough as $|x| \rightarrow \infty$. Let $\widehat{u}(t, \cdot)$ be the Fourier transform of u with respect to the spatial variable: $\widehat{u}(t, \xi) = \int_{\mathbb{R}} u(t, x) e^{-2\pi i \xi x} dx$; the inverse Fourier transform is $u(t, x) = \int_{\mathbb{R}} \widehat{u}(t, \xi) e^{2\pi i \xi x} d\xi$ (assuming that the functions are in appropriate function spaces, so that the Fourier transform and its inverse are well defined; e.g., if $u(t, \cdot) \in \mathcal{S}$ and $\widehat{u}(t, \cdot) \in \mathcal{S}$).

Plug $u(t, \cdot)$ in the heat equation to show that \widehat{G}_t satisfies the ordinary differential equation $\frac{d\widehat{u}(t, \xi)}{dt} = -4\pi^2 \xi^2$, and find the general solution of this equation; this solution will contain one arbitrary function of ξ , let's denote it by $A(\xi)$.

- (b) Now plug u in the initial condition $\lim_{t \rightarrow 0^+} u(t, x) = u_0(x)$, and find the function $A(\xi)$. Write the solution of the initial value problem consisting of the heat equation and the initial condition.

- (c) In the rest of this problem you will obtain an important general form of the solutions of the heat equation. Let $G_t(x)$ stand for the solution of the heat equation corresponding to the choice $A(\xi) \equiv 1$. If $\widehat{G}_t(x) = \int_{\mathbb{R}} G_t(x) e^{-2\pi i \xi x} dx$ is the Fourier transform of G_t (in the spatial variable), then what is \widehat{G}_t ? (This is a trivial question.)

Use the fact that, if $f : \mathbb{R} \rightarrow \mathbb{R}$ is the function given by $f(x) = e^{-\pi a x^2}$ (where $a > 0$ is a constant), then $\widehat{f}(\xi) = \frac{1}{\sqrt{a}} e^{-\pi \xi^2 / a}$ (this is Proposition 8.24 in Folland for $n = 1$), to show that $G_t(x) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}$.

- (d) Check by direct substitution that the function G_t obtained in (c) is a solution of the heat equation.
- (e) An *approximate identity* in \mathbb{R} is a family of functions $\{K_t\}_{t>0}$, where $K_t : \mathbb{R} \rightarrow \mathbb{R}$, with the following properties:

$$\sup_t \int_{\mathbb{R}} |K_t(y)| dy < \infty, \quad \lim_{t \rightarrow 0^+} \int_{\mathbb{R}} K_t(y) dy = 1, \quad \lim_{t \rightarrow 0^+} \int_{\{|y|>\delta\}} |K_t(y)| dy = 0,$$

for any constant $\delta > 0$. (This definition of an approximate identity looks different from the one in Folland, but in fact they are not so different.)

Check that $\{G_t\}_{t>0}$ is an approximate identity. You may use that $\int_{\mathbb{R}} e^{-x^2} dx = \sqrt{\pi}$.

- (f) Prove that the function $u(t, x) := (G_t * u_0)(x)$ satisfies the heat equation and the boundary condition $\lim_{t \rightarrow 0^+} u(t, x) = u_0(x)$, for any u_0 that decays fast enough at infinity. Does one need to impose differentiability conditions on u_0 ? Explain briefly.