## MATH 4163-002 Test II

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March 10, 2016 1:30pm-2:45pm

Answer the questions in the spaces provided on the question sheets. No calculators allowed.

Name: \_

Problem	Score
1	
2	
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Total	

1. (10 points) We say that an operator L is self-adjoint if for any two functions u, v and a multiplication \* on functions, L(u) \* v = u \* L(v). Write down two different examples of self-adjoint operators, if we let \* be standard multiplication.

**Solution:** L(f) = 0 and L(f) = f are the simplest examples.

2. (15 points) Recall that a Fourier series on [-L, L] can be written down in two forms: either as

$$\hat{f}(x) = A_0 + \sum_{n=1}^{\infty} A_n \cos \frac{n\pi x}{L} + B_n \sin \frac{n\pi x}{L},$$

or as

$$\hat{f}(x) = \sum_{m=-\infty}^{\infty} C_m e^{-im\pi x/L}.$$

Express  $C_m$  in terms of the  $A_n$  and  $B_n$  (recall that the *m* in  $C_m$  can be positive or negative, but the *n* in  $A_n, B_n$  cannot be negative.) You may use the trigonometric identities  $2\cos(\theta) = e^{i\theta} + e^{-i\theta}$ , and  $2i\sin(\theta) = e^{i\theta} - e^{-i\theta}$ .

Solution: We plug in the hint to the first version of the Fourier series.

$$\hat{f}(x) = A_0 + \sum_{n=1}^{\infty} A_n \frac{1}{2} \left( \exp \frac{n\pi x}{L} + \exp \frac{-n\pi x}{L} \right) + B_n \frac{1}{2i} \left( \exp \frac{n\pi x}{L} - \exp \frac{-n\pi x}{L} \right),$$

Collecting like terms, we can rewrite this as

$$\hat{f}(x) = A_0 + \sum_{n=1}^{\infty} \frac{A_n + B_n/i}{2} \exp \frac{n\pi x}{L} + \frac{A_n - B_n/i}{2} \exp \frac{-n\pi x}{L}.$$

So clearly, if m is positive,

$$C_m = \frac{A_n - B_n/i}{2}, C_{-m} = \frac{A_n + B_n/i}{2},$$

and

$$C_0 = A_0$$

3. (15 points) Consider a vibrating string lying between x and  $x + \Delta x$ , and let u(x,t) be the height of the string at time t and location x. The vertical acceleration of the string applies a force of  $\rho \Delta x \frac{\partial^2}{\partial t^2} u(x,t)$ . The tension of the string applies a vertical force of  $T \sin(\theta(x + \Delta x, t)) - T \sin(\theta(x, t))$ , where T is the (constant) tension of the string, and  $\theta(x,t)$  is the angle of the string's slope at position x and time t. Also, you may assume that  $\theta(x,t)$  is always small enough that  $\sin(\theta(x,t))$  is always roughly equal to  $\tan(\theta(x,t))$ .

Derive the vibrating string equation from this information (you may want to start by equating all the vertical forces acting on the string)

Solution: We have

$$\rho \Delta x \frac{\partial^2}{\partial t^2} u(x,t) = T \sin(\theta(x + \Delta x, t)) - T \sin(\theta(x,t)).$$

Let us divide by  $\Delta x$  to obtain

$$\rho \frac{\partial^2}{\partial t^2} u(x,t) = \frac{T \sin(\theta(x + \Delta x, t)) - T \sin(\theta(x, t))}{\Delta x}.$$

Taking the limit as  $\Delta x \to 0$  we have

$$\rho \frac{\partial^2}{\partial t^2} u(x,t) = \frac{d}{dx} T \sin(\theta(x,t)).$$

However, notice that the slope of u(x,t) can be expressed in two ways: as  $\frac{\partial}{\partial x}u(x,t)$ , and as rise over run. Rise over run is equivalent to opposite over adjacent, which is equivalent to  $\tan(\theta(x,t))$ , and we can assume that  $\tan(\theta(x,t) \sim \sin(\theta(x,t)))$ . Thus we may assume

$$\frac{\partial}{\partial x}u(x,t) = \sin(\theta(x,t)).$$

Plugging this in to our previous equation, we get

$$\rho \frac{\partial^2}{\partial t^2} u(x,t) = T \frac{\partial^2}{\partial x^2} u(x,t),$$

which is the vibrating string equation.

4. (15 points) Product solutions to the vibrating string equation take the form

$$u_n(x,t) = \sin \frac{n\pi x}{L} \left( A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L} \right), n = 1, 2, 3, \dots$$

Explan what happens to the pitch of the sound produced by the string's vibration if we

- (a) increase the length of the string
- (b) increase the density of the string
- (c) increase the tension of the string

Justify your answers using the above formula for the product solution. Justifications of the form "I know this is true because I play guitar" will earn very little partial credit.

**Solution:** The frequency of the trig functions  $\sin(\omega t)$ ,  $\cos(\omega t)$  is  $\frac{\omega}{2\pi}$  hertz. Thus if we increase L, the  $\omega$  decreases, and the frequency of the vibration decreases, resulting in a lower pitch.

 $c = \sqrt{\frac{T}{\rho}}$ , and so increasing tension increases c and therefore  $\omega$ , leading to a higher pitch. Conversely, increasing the density lowers the T and therefore  $\omega$ , thus decreasing the pitch.

5. (15 points) Consider the Sturm-Liouville operator,

$$L(F(x)) = \frac{d}{dx} \left( p(x) \frac{dF(x)}{dx} \right) + q(x)F(x).$$

You may recall the Green's formula,

$$\int_{a}^{b} (uL(v) - vL(u))dx = p\left(u\frac{dv}{dx} - v\frac{du}{dx}\right)\Big|_{x=a}^{b}$$

Show that if we impose boundary conditions 3f(a) - f'(a) = 0, f'(b) = 0 on both f = u and f = v, the operator L is self-adjoint.

Solution: Self-adjointness is equivalent to showing

$$p\left(u\frac{dv}{dx} - v\frac{du}{dx}\right)\Big|_{x=a}^{b} = 0.$$

We expand out the LHS to get

$$p(b)\left(u(b)\frac{dv(b)}{dx} - v(b)\frac{du(b)}{dx}\right) - p(a)\left(u(a)\frac{dv(a)}{dx} - v(a)\frac{du(a)}{dx}\right).$$

Since u'(b) = v'(b) = 0, the entire first term is zero. This leaves us with

$$p(a)\left(u(a)\frac{dv(a)}{dx}-v(a)\frac{du(a)}{dx}\right).$$

But note that since 3u(a) - u'(a) and 3v(a) - v'(b) are both zero, we must have u'(a) = 3u(a) and v'(a) = 3v(a). This implies

$$p\left(u\frac{dv}{dx} - v\frac{du}{dx}\right)\Big|_{x=a}^{b} = p(a)\left(u(a)3v(a) - v(a)3u(a)\right) = 0.$$

6. (15 points) Let  $\lambda_m$ ,  $\lambda_n$  be two eigenvalues for a self-adjoint Sturm-Liouville equation

$$L(F(x)) + \lambda\sigma(x)F(x) = 0,$$

and let  $F_m(x), F_n(x)$  respectively be their two eigenfunctions. Show that if  $\lambda_m \neq \lambda_n$ , then

$$\int_{a}^{b} F_{m}(x)F_{n}(x)\sigma(x)dx = 0.$$

Hint: what does the fact that L is self-adjoint imply?

## Solution:

Using self-adjointness, we know that

$$\int_{a}^{b} L(F_m)F_n - L(F_n)F_m dx = 0$$

But we also know from the Sturm-Liouville equation that

$$L(F_m(x)) = -\lambda_m \sigma(x) F_m(x),$$

and

$$L(F_n(x)) = -\lambda_n \sigma(x) F_n(x).$$

Plugging these into the self-adjointness equation we get

$$\int_{a}^{b} -\lambda_{m}\sigma(x)F_{m}(x)F_{n}(x) + \lambda_{n}\sigma(x)F_{n}(x)F_{m}(x)dx = 0$$

Factoring, this becomes

$$(\lambda_n - \lambda_m) \int_a^b F_m(x) F_n(x) \sigma(x) dx = 0,$$

but since  $\lambda_n - \lambda_m \neq 0$ , it must be true that

$$\int_{a}^{b} F_{m}(x)F_{n}(x)\sigma(x)dx = 0.$$

7. (15 points) It is known that if our Sturm-Liouville operator L is self-adjoint and if we impose Dirichlet, Neumann, or Robin boundary conditions, then where  $F_1(x), F_2(x)$  are two eigenfunctions corresponding to the same eigenvalue  $\lambda$ ,

$$F_1(x)F_2'(x) - F_2(x)F_1'(x) = 0.$$

Using this, show that  $F_2(x)$  must be a constant multiple of  $F_1(x)$ .

Solution: By the quotient rule,

$$\frac{d}{dx}\left(\frac{F_2(x)}{F_1(x)}\right) = \frac{F_1(x)F_2'(x) - F_2(x)F_1'(x)}{F_1(x)^2} = 0$$

This implies that for some constant C,

$$\frac{F_2(x)}{F_1(x)} = C,$$

or  $F_2(x) = CF_1(x)$ .